

Muskingum-Cunge portraits

revisited

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01 March 2024

ABSTRACT

A comprehensive review of the Muskingum-Cunge method's amplitude and phase portraits is accomplished. Expressions for the amplitude convergence ratio R_1 and phase convergence ratio R_2 are expressed as a function of the following basic numerical parameters: (a) Spatial resolution $L/\Delta x$; (b) Courant number *C*; and (c) weighting factor *X*. An online calculator **Online Muskingum-Cunge Convergence Ratios** is used for convenience.

For practical applications, input may be expressed in terms of relevant hydraulic variables, such as mean velocity, flow depth, channel slope, rating exponent, and flood wave time-of-rise. For this case, an online calculator **Online Muskingum-Cunge Convergence Ratios Practical** is developed.

1. INTRODUCTION

This article revisits the concepts of amplitude and phase portraits, originally due to Leendertse (1967), and their application to the Muskingum-Cunge method of flood routing. An amplitude portrait is a plot of the ratio of numerical to analytical wave amplitudes, i.e., it is a *convergence ratio* with regard to wave amplitude. A phase portrait is a plot of the ratio of numerical to analytical wave phase. Amplitude and phase portraits for the Muskingum-Cunge method were originally developed by **Cunge (1969)** and later expanded for the generalized convection problem by **Ponce** *et al.* **(1979)**.

The portraits are used to assess the stability and convergence properties of a given numerical scheme. This helps determine the range of numerical parameters (spatial resolution, temporal resolution, and their corresponding ratio) for which the scheme may be shown to be stable and convergent. Following **O'Brien** *et al.* (1950), stability is related to round-off errors; convergence is related to truncation errors.

In this article, we revisit the amplitude and phase portraits of the Muskingum-Cunge method with aim to clarify their utility and enable their wider use.

2. MUSKINGUM-CUNGE METHOD

In a seminal paper, Cunge (1969) elucidated the kinematic wave basis of the Muskingum method of flood routing

(1)

(Chow, 1959), leading to the Muskingum-Cunge method (Natural Environment Research Council, 1975; Koussis, 1978; **Ponce and Yevjevich**, **1978**). Cunge's procedure allowed for the calculation of the weighting factor *X* in terms of physical and numerical parameters, instead of relying solely on gaged measurements, as per the original Muskingum method (McCarthy, 1938). This revolutionary concept, based on matching the physical diffusivity of **Hayami (1951)** with the numerical diffusivity of Cunge (1969), enables the Muskingum-Cunge method to remain *grid independent* through a wide range of values of the relevant numerical parameters: spatial resolution, Courant number, and weighting factor. We note that this important feature of grid independence is not shared by any other flood routing method based on the kinematic wave (**Ponce, 1986**).

In his original paper, **Cunge (1969)** expressed the amplitude and phase portraits in terms of the temporal resolution $(T/\Delta t)$ and the reciprocal of the grid ratio $(\Delta x / \Delta t)$. **Ponce** *et al.* **(1979 a)** generalized Cunge's approach to the entire set of numerical schemes that may be construed to solve the general convection problem. They expressed the amplitude and phase portraits in terms of the spatial resolution $(L/\Delta x)$ and the Courant number *C*, the latter defined as the ratio of physical celerity *c* to numerical celerity $(\Delta x / \Delta t)$. In this article, we review the findings of Ponce *et al.* (1979), extensively clarified by **Ponce and Vuppalapati (2016)**, and apply them specifically to the Muskingum-Cunge method.

3. GOVERNING EQUATIONS

The kinematic wave equation is (Ponce, 2014 a):

 $\frac{\partial Q}{\partial t} + c \frac{\partial Q}{\partial x} = 0$

in which Q = discharge, c = kinematic wave celerity, and x and t are space and time variables, respectively.

The discretization of Eq. 1 on the x-t plane (Fig. 1) leads to (Ponce, 2014 b):

$$\frac{X(Q_{j}^{n+1} - Q_{j}^{n}) + (1 - X)(Q_{j+1}^{n+1} - Q_{j+1}^{n})}{\wedge t} + c \frac{(Q_{j+1}^{n} - Q_{j}^{n}) + (Q_{j+1}^{n+1} - Q_{j}^{n+1})}{2 \wedge x} = 0$$
(2)





Simplifying Eq. 2:

$$X(Q_{j}^{n+1} - Q_{j}^{n}) + (1 - X)(Q_{j+1}^{n+1} - Q_{j+1}^{n}) + 0.5C[(Q_{j+1}^{n} - Q_{j}^{n}) + (Q_{j+1}^{n+1} - Q_{j}^{n+1})] = 0$$
(3)

in which C = Courant number, defined as follows:

$$C = c \frac{\Delta t}{\Delta x} \tag{4}$$

The solution of Eq. 3 is postulated in the following exponential form (O'Brien et al., 1950; Cunge, 1969):

$$Q = Q_* e^{i(\sigma x - \beta t)}$$
⁽⁵⁾

in which $\sigma = (2\pi/L)$ is the wavenumber, *L* is the wavelength of the disturbance, and $\beta = (2\pi/T)$ is the wave period (**Ponce and Simons, 1977**).

4. CONVERGENCE RATIOS

The derivation of the convergence ratios has been described in detail by **Ponce and Vuppalapati (2016)**. The convergence ratio with regard to wave amplitude is:

$$R_{1} = \frac{[(P - M)^{2} + N^{2}]^{1/2}}{P}$$
(6)

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in which:	
P = 1 - 2 (1 - p)(1 - S) S	(6a)
M = (1-p)(1-2S)C	(6b)
N = -Cq	(6c)
wherein	
$p = \cos(\sigma \Delta x)$	(6d)
$q = \sin(\sigma \Delta x).$	(6e)
S = X - (C/2)	(6f)
and $C = \text{Courant number (Eq. 4)}.$	

The convergence ratio with regard to wave phase is:

$$R_2 = \frac{1}{C \sigma \Delta x} \tan^{-1} \left(\frac{N}{P - M} \right)$$
(7)

Therefore, given: (1) weighting factor X (in Eq. 2), (2) spatial resolution $L/\Delta x$ [from which $\sigma \Delta x = (2\pi / L)\Delta x$ = $2\pi / (L/\Delta x)$], and (3) Courant number C (Eq. 4), Eqs. 6 and 7 enable the calculation of the Muskingum-Cunge amplitude and phase convergence ratios, respectively. The calculation of the convergence ratios is accomplished using the calculator **Online Muskingum-Cunge Convergence Ratios**.

5. AMPLITUDE AND PHASE PORTRAITS

The plotting of convergence ratios R_1 and R_2 as a function of Courant number *C* (in the abscissas) and spatial resolution $L/\Delta x$ (as the curve parameter) leads to a set of amplitude and phase portraits R_1 and R_2 , respectively, a pair for each value of weighting factor *X* in the range $0.0 \le X \le 0.5$. We observe that this range of analysis for the weighting factor *X* is amply justified by kinematic wave theory (**Ponce** *et al.*, **1979 a**).

Figures 2 to 8 show the amplitude and phase portraits of the Muskingum-Cunge method (Eq. 3), for the following conditions: (a) Courant number $0.1 \le C \le 4.0$, (b) spatial resolution $20 \le L/\Delta x \le 100$, shown at intervals of $L/\Delta x = 20$, and (c) weighting factor $0.0 \le X \le 0.6$, shown at intervals of X = 0.1 (7 × 2 = 14 graphs). Note that the value X = 0.6 (Fig. 8) is included here only for comparison purposes, i.e., to show the numerical effect of a value X > 0.5, which results in numerical amplification, R > 1; see Fig. 8 (a).

The selected range of Courant numbers shown ($0.1 \le C \le 4.0$) encompasses the central value C = 1, for which absolute accuracy is obtained for X = 0.5; see Figs. 7 (a) for amplitude portrait and 7 (b) for phase portrait. The range of spatial resolution shown in Figures 2 to 8 ($20 \le L/\Delta x \le 100$) depicts five (5) values, varying from low ($L/\Delta x = 20$; poor resolution) to high ($L/\Delta x = 100$; very good resolution). The spatial resolution is the number of discrete intervals Δx per wavelength *L*.







Fig. 3 (a) Amplitude portrait for X = 0.1.









Fig. 3 (b) Phase portrait for X = 0.1.









1

0.96

0.94

0.92

0.9

0.88

0.86





Fig. 6 (a) Amplitude portrait for X = 0.4.





Fig. 6 (b) Phase portrait for X = 0.4.







6. ANALYSIS

The examination of Figs. 2 to 8 leads to the following conclusions:

- 1. Perfect convergence amounts to $R_1 = 1$ and $R_2 = 1$, i.e., implying that Eq. 2 is the exact solution of Eq. 1. It is achieved only for X = 0.5 [confirm with Fig. 7 (a)] and Courant number C = 1 [confirm with Fig. 7 (b)]. This behavior is due to the following reasons: (1) for X = 0.5, the numerical scheme (Eq. 2) is fully centered (Fig. 1) and, therefore, the *second-order error* vanishes; and (2) for C = 1, the *third-order error* also vanishes (**Ponce** *et al.*, 1979 b).
- For any *X*-*C* pair other than *X* = 0.5 and *C* = 1, the numerical solution of Eq. 2 is *only* an approximation of the analytical solution of Eq. 1. The accuracy of the approximation is a function of the following parameters:
 (a) spatial resolution *L*/Δ*x*; (b) Courant number *C*; and weighting factor *X* of the numerical scheme (Fig. 1).
- 3. As expected, convergence in both amplitude and phase improves with an increase in spatial resolution. In Figures 2 to 8, the low resolution ($L/\Delta x = 20$) is depicted in red color, while the high resolution is shown in black. In particular, note the marked departure of the red-colored curves from the *X*-axis (R = 1), indicating a substantial lack of convergence for the low resolution case.
- 4. Convergence in both amplitude and phase degrades with an increase in Courant number as *C* increases above 1, particularly as $C \gg 1$; for example, for C = 4.
- 5. Convergence in amplitude improves *substantially* with an increase of X in the range $0 \le X < 0.5$. At X = 0.5, amplitude convergence is perfect for *all* Courant numbers [see Fig. 7 (a)].
- 6. Convergence in phase improves *slightly* with an increase of X in the range $0 \le X < 0.5$. At X = 0.5, phase convergence is perfect only for C = 1 [see Fig. 7 (b)].
- 7. At low spatial resolutions, refer to red curves in Figs. 2 (b) to 8 (b), convergence in phase (as depicted by R_2) degrades significantly as Courant numbers decrease below C = 1, leading to $R_2 > 1$, i.e., numerical instability [see Figs. 5 (b), 6 (b), y 7(b)].
- 8. For X > 0.5, as shown in Figs. 8(a) and 8(b), i.e., *outside* of the feasible range of X in Muskingum-Cunge routing, the amplitude convergence ratio R_1 becomes positive (amplification) and, accordingly, stability and

convergence degrade in an explosive manner.

We conclude that the Muskingum-Cunge numerical model is a good representation of the analytical prototype, provided:

- 1. The spatial resolution $L/\Delta x$ is sufficiently high, preferably of the order of $L/\Delta x \approx 100$.
- 2. The Courant number C is sufficiently close to 1.
- 3. The weighting factor X is high enough in the range $0.0 \le X < 0.5$; preferably between 0.3 and 0.5.

Based on the analysis presented herein, the recommended parameter values for reasonably adequate convergence are the following: (a) spatial resolution $L/\Delta x \approx 100$; (b) Courant number $C \approx 1$; and (c) weighting factor X between 0.3 and 0.5.

7. APPLICATIONS

In practice, the convergence of the Muskingum-Cunge method depends on the values of Courant and cell Reynolds numbers, which in turn depend on the hydraulics of the flow and the chosen grid size (Δx and Δt) (**Ponce, 2014 b**). A set of amplitude and phase convergence ratios may be calculated in terms of hydraulic and hydrologic variables.

To this end, define the time-of-rise t_r of the flood wave as the time it takes it to reach the peak flow. Assuming a sinusoidal perturbation as a first approximation, the time base, or wave period *T* is:

$$T = 2 t_r \tag{8}$$

The wave celerity c is:

$$c = \beta u$$

in which *u* is the mean flow velocity.

In turbulent open-channel flow, the value of β varies from $\beta = 5/3$ for a hydraulically wide channel, down to $\beta = 1$ for an inherently stable channel (**Ponce and Porras, 1995**; **Ponce, 2014 c**). For a rectangular channel, β may be approximated as 5/3, for a triangular channel as 4/3, and for a hydraulically wide natural channel, a reasonably good approximation is $\beta = 1.6$.

The wavelength *L* of the perturbation is:

$$L = c T \tag{10}$$

The unit-width discharge q is:

$$q = u d$$

in which *d* is the mean flow depth. The spatial resolution is $L/\Delta x$.

The Courant number (Eq. 4), repeated here for convenience with a new equation number, is:

(11)

(9)

$$C = c \frac{\Delta t}{\Delta x}$$
(12)

The cell Reynolds number is (Ponce, 2014e):

$$D = \frac{q}{S_o c \Delta x}$$
(13)

By definition, the Muskingum-Cunge weighting factor *X* is:

$$X = \frac{1}{2} (1 - D)$$
(14)

For a given case, Eqs. 6 and 7, with support from Eqs. 8 to 14, enable the calculation of the convergence ratios of the Muskingum-Cunge method. Two cases are explained here: (1) theoretical, and (2) practical.

The calculation of the theoretical Muskingum-Cunge convergence ratios may be accomplished using the calculator **Online Muskingum-Cunge Convergence Ratios**. Three examples are discussed here: (**A**) perfect convergence, (**B**) good convergence, and (**C**) poor convergence. For Case **A**, assume X = 0.5; $L /\Delta x = 100$; and C = 1. The results, shown in Fig. 9, are: $R_1 = 1$, and $R_2 = 1$, confirming the perfect convergence of the Muskingum-Cunge method for the case of X = 0.5, $L /\Delta x = 100$, and C = 1.



Fig. 9 Muskingum-Cunge convergence ratios for X = 0.5, $L/\Delta x = 100$, and C = 1: $R_1 = 1$; $R_2 = 1$.

For Case **B**, assume X = 0.35; $L/\Delta x = 60$; and C = 1.1. The results, shown in Fig. 10, are: $R_1 = 0.998195$; $R_2 = 0.999563$, confirming the good convergence of the Muskingum-Cunge method for the case of X = 0.35, $L/\Delta x = 60$, and C = 1.1.



Fig. 10 Muskingum-Cunge convergence ratios for X = 0.35, $L/\Delta x = 60$, and C = 1.1: $R_1 = 0.998185$; $R_2 = 0.999563$.

For Case **C**, assume X = 0.05; $L /\Delta x = 20$; and C = 2.5. The results, shown in Fig. 11, are: $R_1 = 0.908286$; $R_2 = 0.944956$, confirming the *poor* convergence (i.e., *R* ratios much smaller than 1) of the Muskingum-Cunge method for the case of X = 0.05, $L /\Delta x = 20$, and C = 2.5.



Fig. 11 Muskingum-Cunge convergence ratios for X = 0.05, $L/\Delta x = 20$, and C = 2.5: $R_1 = 0.908286$; $R_2 = 0.944956$.

The calculation of the practical Muskingum-Cunge convergence ratios may be accomplished using the calculator **Online Muskingum-Cunge Convergence Ratios Practical**. Two examples are discussed here: (**D**) small stream, and (**E**) large stream. For Example **D**, assume the following data set:

- 1. Mean velocity u = 1.75 m/s;
- 2. Mean flow depth d = 2.3 m;
- 3. Mean channel slope S = 0.0013;
- 4. Rating exponent $\beta = 1.6$;
- 5. Flood wave time-of-rise $t_r = 24$ hr;
- 6. Reach length $\Delta x = 4800$ m; and
- 7. Time interval $\Delta t = 0.5$ hr.

The application of **Online Muskingum-Cunge Convergence Ratios Practical**, shown in Fig. 12, leads to: $R_1 = 0.99953$; $R_2 = 0.999915$. In this case, the spatial resolution is $L/\Delta x = 100.8$; the Courant number is C = 1.05; and the weighting factor is X = 0.385. All three parameters are confirmed to lie within the recommended ranges for adequate convergence.



Fig. 12 Muskingum-Cunge convergence ratios for Example **D**: $R_1 = 0.99953$; $R_2 = 0.999915$.

For the next, and last, example E, assume the following data set:

- 1. Mean velocity u = 0.5 m/s;
- 2. Mean flow depth d = 6.2 m;
- 3. Mean channel slope S = 0.018;
- 4. Rating exponent $\beta = 1.65$;
- 5. Flood wave time-of-rise $t_r = 36$ hr;
- 6. Reach length $\Delta x = 2500$ m; and
- 7. Time interval $\Delta t = 1.5$ hr.

The application of **Online Muskingum-Cunge Convergence Ratios Practical**, shown in Fig. 13, leads to: $R_1 = 0.9996$; $R_2 = 0.999014$. In this case, the spatial resolution is $L/\Delta x = 85.536$; the Courant number is C = 1.728; and the weighting factor is X = 0.458. All three parameters are confirmed to lie within the recommended ranges for adequate convergence.



Fig. 13 Muskingum-Cunge convergence ratios for Example E: $R_1 = 0.9996$; $R_2 = 0.999014$.

8. CONCLUDING REMARKS

Herein we endeavor to clarify some of the concepts elaborated in this article, with aim to encourage its wider use.

On the Courant number. Numerical convergence, which may be taken as numerical accuracy, depends on spatial resolution $L/\Delta x$ and temporal resolution $T/\Delta t$, both of these being dimensionless expressions of grid size.

In computacional hydraulics, the Courant number *C* is defined as the ratio of physical celerity to grid "celerity": $C = c / (\Delta x / \Delta t) = c (\Delta t / \Delta x) = (L/T) (\Delta t / \Delta x) = (L/\Delta x) / (T/\Delta t)$. In other words, the Courant number is indeed the ratio of spatial to temporal resolution. Only two of the three variables: (1) spatial resolution, (2) temporal resolution, and (3) Courant number, are required to fully define the analysis. As stated in Section 2 of this article, we use spatial resolution and Courant number, as shown in Figs. 2 to 8.

On the limitations of the Muskingum method. The classical Muskingum method of flood routing (Chow, 1959) solves the kinematic wave equation numerically, thus, introducing a certain undefined amount of numerical diffusion. In practice, the latter manifests itself as an appreciable amount of attenuation, or diffusion, of the flood wave being calculated. There is no relation of this introduced numerical diffusion with the true physical diffusion, if any, of the problem under consideration [*Note:* Recall that the true solution of a kinematic wave does not allow for *any* physical diffusion]. Therefore, the classical Muskingum method fails to solve the flood propagation problem in a meaningful way. The actual solution will be heavily dependent on the chosen grid size, which has not been explicitly related to the physical problem at hand.

On the strength of the Muskingum-Cunge method. Unlike the Muskingum method, the Muskingum-Cunge method (M-C method) is physically and numerically based. The weighting factor X is calculated by Eq. 14, which in turn is based on physical and numerical variables (Eq. 13). This remarkable feature of the M-C method is unique in the field of computational hydraulics, enabling it to become grid independent through a wide range of values of spatial resolution. Once the spatial grid interval Δx is selected or determined a priori, the weighting factor X is fixed by Eq. 14. In this way, the values of physical and numerical diffusion have been matched, and this procedure renders the method and its result grid independent!

On the meaning of the amplitude and phase portraits. Since the kinematic wave equation (Eq. 1) exhibits no physical diffusion (i.e., damping), Eq. 6 (amplitude portrait) is a measure of the numerical diffusion (i.e., a value less than 1) or, alternatively, numerical amplification (i.e., *negative diffusion*, a value greater than 1), after one time increment Δt (Ponce *et al.*, 1979 d). Correspondingly, a similar conclusion for wave phase (either retardation or acceleration) follows from Eq. 7 (phase portrait). It should not escape our attention that the closeness to 1 of both portrait values shown in Figs. 9 to 13 reflects the fact that the values are applicable per time increment Δt .

9. SUMMARY

A comprehensive review of the Muskingum-Cunge method's amplitude and phase portraits is accomplished. Expressions for the amplitude convergence ratio R_1 and phase convergence ratio R_2 are expressed as a function of the following basic numerical parameters: (a) spatial resolution $L/\Delta x$; (b) Courant number *C*; and (c) weighting factor *X*. An online calculator **Online Muskingum-Cunge Convergence Ratios** is used for convenience.

For practical applications, input may be expressed in terms of relevant hydraulic variables, such as mean velocity, flow depth, channel slope, rating exponent, and flood wave time-of-rise. For this case, an online calculator **Online Muskingum-Cunge Convergence Ratios Practical** is developed.

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